

QUANTUM COMPUTING

MS 2025 - Sheffield

Venanzio Capretta

LECTURE 3 : Multiple qubit states / gates

The no-cloning theorem

Quantum Teleportation

MULTIPLE QUBIT STATES

The state space for 1 qubit is \mathbb{C}^2
The state space for n qubits is:

$$\underbrace{\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2}_{n \text{ times}} \cong \mathbb{C}^{2^n}$$

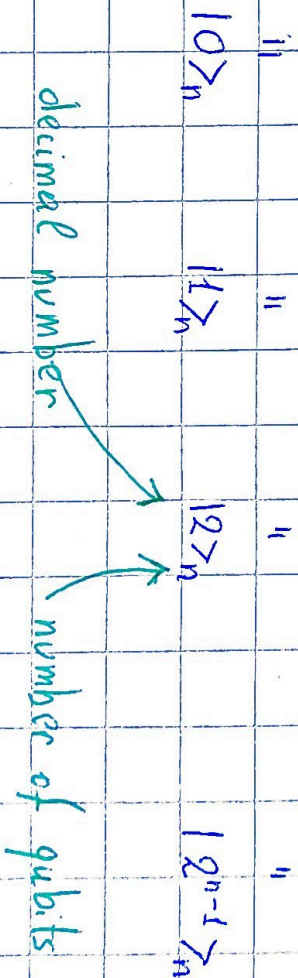
← tensor product

It's the space generated by the tensor products:

$$(\alpha_1 |0\rangle + \beta_1 |1\rangle) \otimes (\alpha_2 |0\rangle + \beta_2 |1\rangle) \otimes \dots \otimes (\alpha_n |0\rangle + \beta_n |1\rangle)$$

If a state can be written in this form, it's called separable
Otherwise it is called entangled

Computational Basis: $|00\dots 0\rangle, |0\dots 01\rangle, |0\dots 10\rangle, \dots, |1\dots 11\rangle$



BELL STATES

maximally entangled states on 2 qubits:

$$|\Phi^+\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

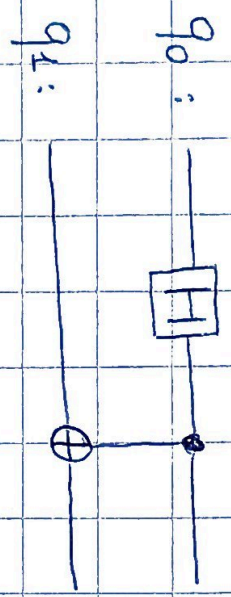
$$|\Psi^+\rangle = \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle)$$

$$|\Phi^-\rangle = \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$$

$$|\Psi^-\rangle = \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle)$$

$\{|\Phi^+\rangle, |\Phi^-\rangle, |\Psi^+\rangle, |\Psi^-\rangle\}$ is an orthonormal basis of $\mathbb{C}^2 \otimes \mathbb{C}^2$

Circuit to create a Bell state:



$$|00\rangle \xrightarrow{H} \frac{1}{\sqrt{2}} (|00\rangle + |10\rangle) \xrightarrow{\text{CNOT}} \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

$$|\Phi^+\rangle$$

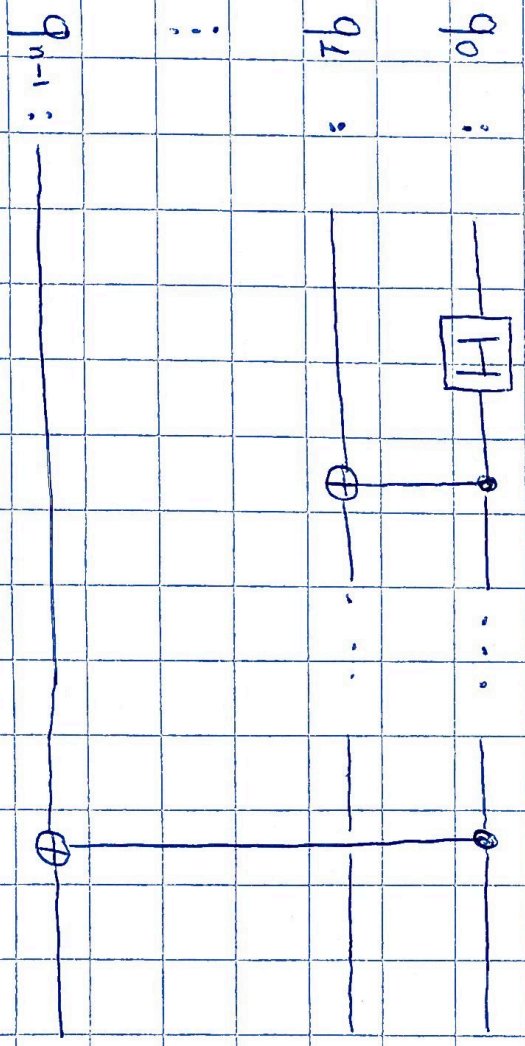
Exercises: draw circuits to create the other Bell states

GENERALIZATION TO n QUBITS:

Greenberger - Horne - Zeilinger (GHZ) states:

$$|GHZ\rangle_n = \frac{1}{\sqrt{2}} (|0\rangle^{\otimes n} + |1\rangle^{\otimes n}) = \frac{1}{\sqrt{2}} (|0\rangle_n + |2^n-1\rangle_n)$$

Created with the GHZ circuit:



THE QUANTUM $H^{\otimes n}$ GATES

Apply the H gate to each qubit in parallel.

For 2 qubits: $H^{\otimes 2} (|Y_1\rangle \otimes |Y_2\rangle) = (H|Y_1\rangle) \otimes (H|Y_2\rangle)$

How to compute tensor products of states and operators:

States:

$$|Y_1\rangle = \alpha_1 |0\rangle + \beta_1 |1\rangle$$

$$|Y_2\rangle = \alpha_2 |0\rangle + \beta_2 |1\rangle$$

$$|Y_1\rangle \otimes |Y_2\rangle = \alpha_1 \alpha_2 |00\rangle + \alpha_1 \beta_2 |01\rangle + \beta_1 \alpha_2 |10\rangle + \beta_1 \beta_2 |11\rangle$$

Operators:

$$A = \begin{bmatrix} a_{00} & \dots & a_{0n} \\ \vdots & \ddots & \vdots \\ a_{no} & \dots & a_{nn} \end{bmatrix}$$

$$B = \begin{bmatrix} b_{00} & \dots & b_{0m} \\ \vdots & \ddots & \vdots \\ b_{mo} & \dots & b_{mm} \end{bmatrix}$$

$$A \otimes B = \begin{bmatrix} a_{00}B & \dots & a_{0n}B \\ \vdots & \ddots & \vdots \\ a_{no}B & \dots & a_{nn}B \end{bmatrix} =$$

$$\begin{bmatrix} a_{00}b_{00} & \dots & a_{00}b_{0m} & \dots & a_{0n}b_{00} & \dots & a_{0n}b_{0m} \\ \vdots & \ddots & \vdots & \ddots & \vdots & \ddots & \vdots \\ a_{no}b_{00} & \dots & a_{no}b_{0m} & \dots & a_{no}b_{no} & \dots & a_{no}b_{nm} \end{bmatrix}$$

So, For the Hadamard Gate:

$$H^{\otimes 2} = H \otimes H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \otimes H = \frac{1}{\sqrt{2}} \begin{bmatrix} H & H \\ H & -H \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & -1 & 1 \end{bmatrix}$$

Recursively:

$$H^{\otimes (n+1)} = H \otimes H^{\otimes n} = \frac{1}{\sqrt{2}} \begin{bmatrix} H^{\otimes n} & H^{\otimes n} \\ H^{\otimes n} & -H^{\otimes n} \end{bmatrix}$$

$H^{\otimes n}$ is often used at the beginning of a circuit to prepare the qubits in a state of maximum superposition:

$$H^{\otimes 2} |00\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |01\rangle + |10\rangle + |11\rangle) \quad H^{\otimes n} |0\rangle_n = \frac{1}{\sqrt{2^n}} (|0\rangle_n + \dots + |2^n - 1\rangle_n)$$

General form of an n-qubit state:

$$|\psi\rangle = \sum_{j=0}^{2^n-1} \alpha_j |j\rangle_n \quad \text{where } \sum_{j=0}^{2^n-1} |\alpha_j|^2 = 1$$

So $H^{\otimes n} |0\rangle_n = \frac{1}{\sqrt{2^n}} \sum_{j=0}^{2^n-1} |j\rangle_n$

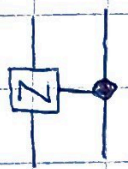
CONTROLLED GATES

CNOT gate



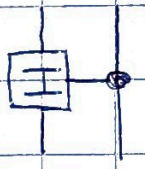
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

CZ gate



$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

Controlled Hadamard gate: CH



$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1/\sqrt{2} & 1/\sqrt{2} \\ 0 & 0 & 1/\sqrt{2} & -1/\sqrt{2} \end{bmatrix}$$

Generalized controlled gate: If gate U has matrix

$$\begin{bmatrix} U_{00} & U_{01} \\ U_{10} & U_{11} \end{bmatrix}$$

Then the controlled U gate has matrix

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & U_{00} & U_{01} \\ 0 & 0 & U_{10} & U_{11} \end{bmatrix}$$

Circuit equivalences (used to simplify circuits)

(General way to reason about quantum diagrams: ZX calculus)

$$\boxed{Z} = \boxed{H} - \oplus - \boxed{H}$$

$$\text{---} \oplus \text{---} = \boxed{H} - \boxed{Z} - \boxed{H}$$

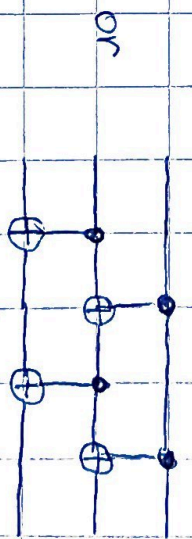
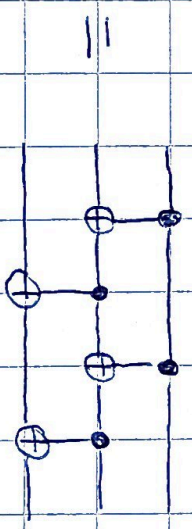
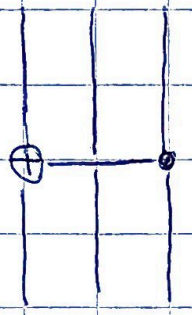
SWAP:



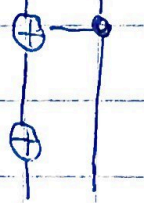
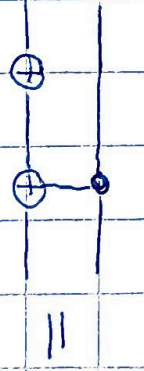
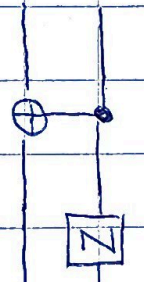
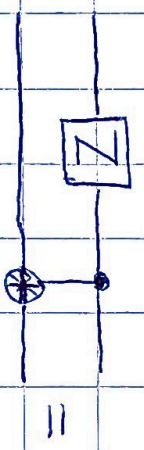
inverted CNOT



Bridge operations:



Commutation Rules:



NO-CLONING THEOREM

Quantum Information Cannot Be Copied

Assume there is a gate CLONE that can copy the state of a qubit



$$\text{If } |\psi\rangle = \alpha|10\rangle + \beta|11\rangle$$

$$\text{Then } \text{CLONE } |\psi\rangle = |\psi\psi\rangle$$

$$\text{Left-hand side: } \text{CLONE} \left((\alpha|10\rangle + \beta|11\rangle) \otimes |10\rangle \right) = \text{CLONE} (\alpha|100\rangle + \beta|110\rangle)$$

$$\stackrel{\text{Linearity}}{=} \alpha \cdot \text{CLONE } |100\rangle + \beta \cdot \text{CLONE } |110\rangle = \boxed{\alpha|100\rangle + \beta|111\rangle}$$

$$\text{Right-hand side: } |\psi\psi\rangle = (\alpha|10\rangle + \beta|11\rangle) \otimes (\alpha|10\rangle + \beta|11\rangle)$$

$$= \alpha^2|100\rangle + \alpha\beta|101\rangle + \beta\alpha|110\rangle + \beta^2|111\rangle$$

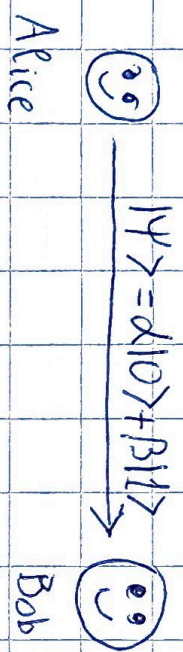
The two sides are not generally equal (eg if $\alpha\beta \neq 0$)

QUANTUM TELEPORTATION

Transfer a qubit, preserving state

We can transfer a qubit perfectly using:

3 qubits + 2 classical bits



We need:

- 2 entangled bits each party has half of the entangled pair
- The message qubit to be teletransported
- A classical communication line between the two parties to transmit 2 classical bits
(so this operation cannot be faster than the speed of light)

STEP 1 Create an entangled pair of qubits and distribute it to both parties



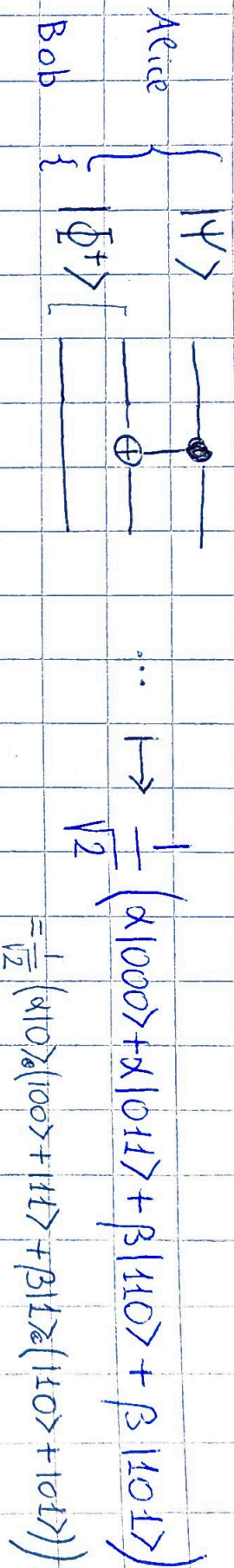
(This is done in advance, maybe before Alice and Bob separate)

Alice also has the message qubit $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$

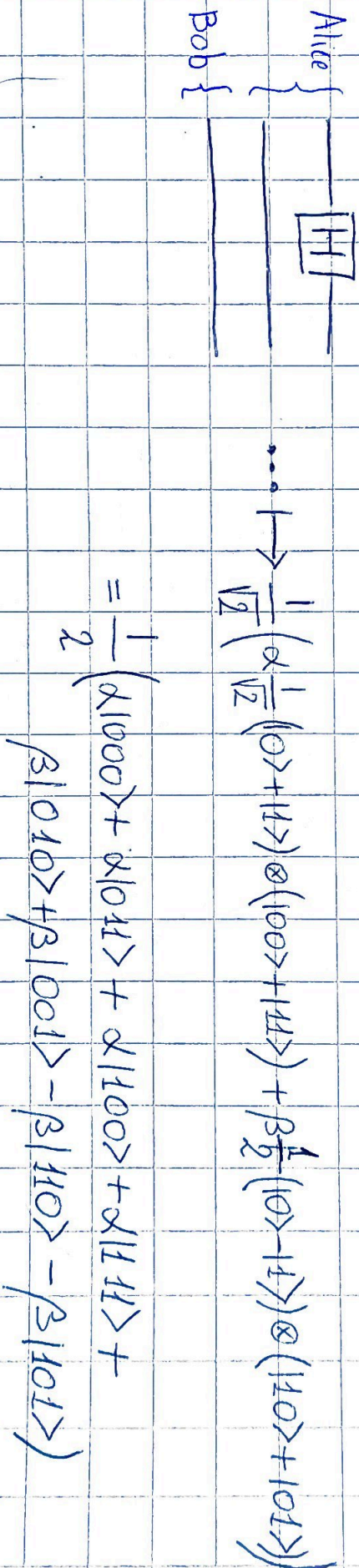
Complete quantum state:

$$|\psi\rangle \otimes |\Phi^+\rangle = \frac{1}{\sqrt{2}} (\alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle)$$

STEP 2 Alice applies a CNOT to her qubits:

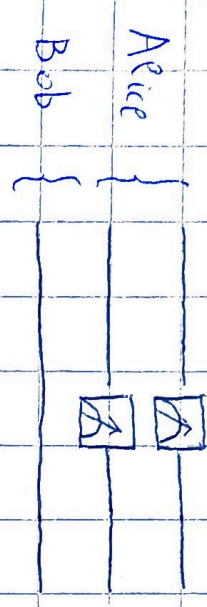


STEP 3 Alice applies a Hadamard gate to the message qubit:



$$= \frac{1}{2} \left(|00\rangle \otimes (\alpha|10\rangle + \beta|11\rangle) + |01\rangle \otimes (\alpha|11\rangle + \beta|10\rangle) + |10\rangle \otimes (\alpha|10\rangle - \beta|11\rangle) + |11\rangle \otimes (\alpha|11\rangle - \beta|10\rangle) \right)$$

STEP 4 Alice measures her qubits :



There are 4 possible results of the measurement:

$|00\rangle, |01\rangle, |10\rangle, |11\rangle$
each with probability $1/4$

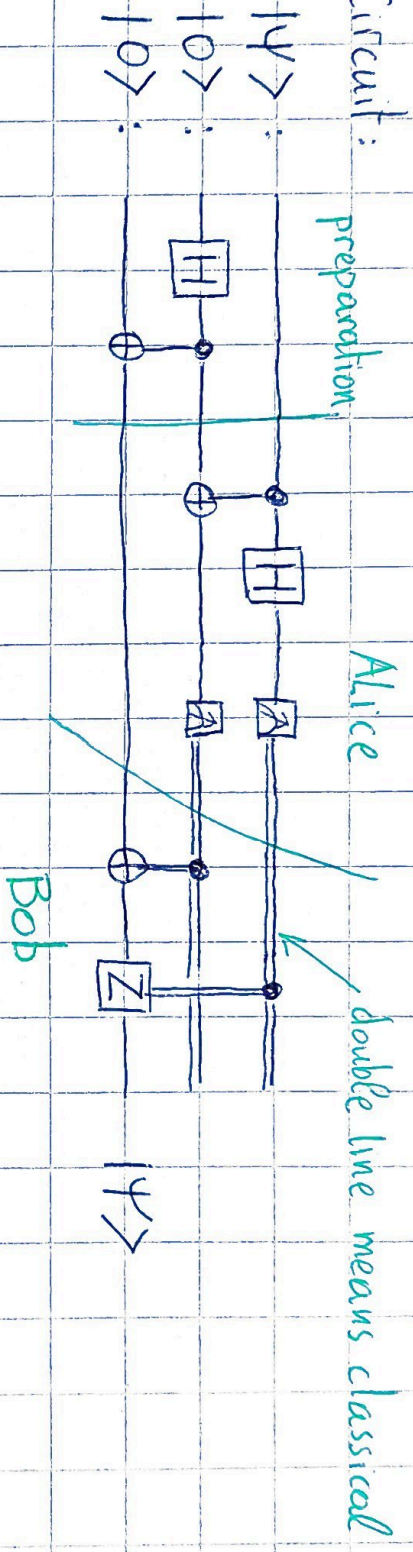
- $|00\rangle \rightarrow$ new state $|00\rangle \otimes (\alpha|10\rangle + \beta|11\rangle)$
- $|01\rangle \rightarrow$ new state $|01\rangle \otimes (\alpha|11\rangle + \beta|10\rangle)$
- $|10\rangle \rightarrow$ new state $|10\rangle \otimes (\alpha|10\rangle - \beta|11\rangle)$
- $|11\rangle \rightarrow$ new state $|11\rangle \otimes (\alpha|11\rangle - \beta|10\rangle)$

STEP 5 Alice transmits the result of the measurement along the classical channel

STEP 6 Bob applies to his qubit :- a Z gate if the first bit is 1
- a NOT gate if the second bit is 1

This sets Bob's qubit to $|1\rangle$

FuPe Circuit:



Advantages of teleportation : • We can move a quantum state with just classical communication

• It's easier to move known qubit states (the two entangled qubits) than unknown ones (the message qubit)

Uses: • Reducing errors

• Linking quantum computers into networks

• Ultra-secure communication channels